## **Two-Dimensional Wave Equations**

Swiss Mathematician and physicist Leonhard Euler discovered the wave equation in three space dimensions.

Leonhard Euler (1707 – 1783) was a pioneering Swiss mathematician and physicist. He made important discoveries in fields as diverse as infinitesimal calculus and graph theory. He is also renowned for his work in mechanics, fluid dynamics, optics, astronomy, and music theory



# TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations**

- The wave equation is an important second-order linear partial differential equation for the description of waves – as they occur in physics – such as sound waves, light waves and water waves.
- It arises in fields like acoustics, electromagnetics, and fluid dynamics.



A solution of the wave equation in two dimensions with a zero-displacement boundary condition along the entire outer edge.

# **Two-Dimensional Wave Equations**

Consider the motion of a rectangular membrane (in the absence of gravity) using the two-dimensional wave equation. Plots of the spatial part for modes are illustrated below.  $\psi_{11}$   $\psi_{12}$   $\psi_{13}$ 



# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

Consider the motion of a rectangular membrane (in the absence of gravity) using the two-dimensional wave equation.



# **Two-Dimensional Wave Equations**

The governing equations of motion that describe the propagation of waves in situations involving two independent variables appear typically as

$$\nabla \cdot (k \nabla u(x, y, t)) - \rho \frac{\partial^2 u}{\partial t^2} + f(x, y, t) = 0$$
 in  $\Omega$ 

$$u = g(s, t)$$
 on  $\Gamma_1$ 

$$k \frac{\partial u}{\partial n} + \alpha(s,t) u = q(s,t)$$
 on  $\Gamma_2$ 

$$u(x, y, 0) = c(x, y)$$
 in  $\Omega$ 

$$\frac{\partial u(x,y,0)}{\partial t} = d(x,y) \qquad \text{in } \Omega$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

Where  $\Omega$  is the interior domain, and  $\Gamma_{\rm 1}$  and  $\Gamma_{\rm 2}$  form the boundary of the domain.



## **Two-Dimensional Wave Equations**

The physical constants are *k* and  $\rho$ , and *s* is a linear measure of the position on the boundary.



#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

The type of boundary condition specified on  $\Gamma_2$  results from a local balance between internal and external forces.



## **Two-Dimensional Wave Equations**

- This equation is similar in form to the parabolic initialboundary value problem presented in the previous section with the very important change in the time derivative term from a first to a second derivative, and with the addition of a second initial condition on the velocity.
- With these changes, the problem changes from a parabolic initial-boundary value problem to a hyperbolic initial boundary value problem.

# TIME-DEPENDENT PROBLEMS

- The basic steps of discretization, interpolation, elemental formulation, assembly, constraints, solution, and computation of derived variables are presented in this section as they relate to the two-dimensional hyperbolic initial-boundary value problem.
- The Galerkin method, in connection with the corresponding weak formulation to be developed, will be used to generate the finite element model.
- **Discretization** Referred to the material in Chapter 3.

#### **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

- **Interpolation -** The solution is assumed to be expressible in terms of the nodally based interpolation functions  $n_i(x, y)$  introduced and discussed in Section 3.2.
- In the present setting, these interpolation functions are used with the semidiscretization.

$$u(x, y, t) = \sum_{i=1}^{N+1} u_i(t) n_i(x, y)$$

The  $n_i(x, y)$  are nodally based interpolation functions and can be linear, quadratic, or as otherwise desired.

#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations**

#### The Galerkin Finite Element Method

- **Elemental Formulation -** The starting point for the elemental formulation is the weak formulation of the initial-boundary value problem.
- The first step in developing the weak formulation is to multiply the differential equation by an arbitrary test function v(x, y) vanishing on  $\Gamma_1$ .

The result is then integrated over the domain  $\Omega$  to obtain

$$\iint_{\Omega} v \left( \nabla \cdot \left( k \nabla u \right) - \rho \frac{\partial^2 u}{\partial t^2} + f \right) d\Omega = 0$$

#### **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

**Elemental Formulation -** Using the two-dimensional form of the divergence theorem to integrate the first term by parts, the results after rearranging are:

$$\iint_{\Omega} \nabla v \cdot k \nabla u \, d\Omega + \iint_{\Omega} v \rho \frac{\partial^2 u}{\partial t^2} \, d\Omega$$
$$= \int_{\Gamma} v \mathbf{n} \cdot k \nabla u \, d\Gamma + \iint_{\Omega} v f \, d\Omega$$

where  $\Gamma = \Gamma_1 + \Gamma_2$ 

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

**Elemental Formulation** - Recalling that *v* vanishes on  $\Gamma_1$ and that  $k\partial u/\partial x = k\mathbf{n} \cdot \nabla u = q - hu$  on  $\Gamma_2$ , it follows that

$$\iint_{\Omega} \nabla v \cdot k \nabla u \, d\Omega + \iint_{\Omega} v \rho \frac{\partial^2 u}{\partial t^2} \, d\Omega$$
$$= \int_{\Gamma} v (q - hu) \, d\Gamma + \iint_{\Omega} v f \, d\Omega$$

This equation is the required weak formulation for the twodimensional diffusion problem.

# **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

Elemental Formulation - Substituting the approximation of

u(x, y, t) into the weak formulation and taking  $v = n_{k'}$ k = 1, 2, ... yields:

$$\iint_{\Omega^{-}} \left( \frac{\partial n_{k}}{\partial x} k \sum u_{i} \frac{\partial n_{i}}{\partial x} + \frac{\partial n_{k}}{\partial y} k \sum u_{i} \frac{\partial n_{i}}{\partial y} \right) d\Omega$$
$$+ \iint_{\Omega^{-}} n_{k} \rho \sum \ddot{u}_{i} n_{i} d\Omega + \int_{\Gamma^{-}} n_{k} h \sum u_{i} n_{i} d\Gamma$$
$$= \iint_{\Omega^{-}} n_{k} f d\Omega + \int_{\Gamma^{-}_{2}} n_{k} q d\Gamma \qquad k = 1, 2, ...$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

**Elemental Formulation –** Where  $\Omega^{-}$  and  $\Gamma_{2}^{-}$  represent the elemental areas approximating  $\Omega$ , and the collection of the elemental edges approximating  $\Gamma_{2}$ , respectively.

$$\iint_{\Omega^{-}} \left( \frac{\partial n_{k}}{\partial x} k \sum u_{i} \frac{\partial n_{i}}{\partial x} + \frac{\partial n_{k}}{\partial y} k \sum u_{i} \frac{\partial n_{i}}{\partial y} \right) d\Omega$$
$$+ \iint_{\Omega^{-}} n_{k} \rho \sum \ddot{u}_{i} n_{i} d\Omega + \int_{\Gamma^{-}} n_{k} h \sum u_{i} n_{i} d\Gamma$$
$$= \iint_{\Omega^{-}} n_{k} f d\Omega + \int_{\Gamma^{-}_{2}} n_{k} q d\Gamma \qquad k = 1, 2, ...$$

# **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

**Elemental Formulation –** This *N* x *N* set of linear algebraic equations can be written as

$$\sum_{i=1}^{N} \left( A_{ki} u_{i} + B_{ki} \ddot{u}_{i} \right) = F_{k}(t) \qquad k = 1, 2, ..., N$$
$$A_{ki} = \iint_{\Omega^{-}} \left( \frac{\partial n_{k}}{\partial x} k \frac{\partial n_{i}}{\partial x} + \frac{\partial n_{k}}{\partial y} k \frac{\partial n_{i}}{\partial y} \right) d\Omega + \int_{\Gamma^{-}} n_{k} h n_{i} d\Gamma$$
$$B_{ki} = \iint_{\Omega^{-}} n_{k} \rho n_{i} d\Omega$$
$$F_{k} = \iint_{\Omega^{-}} n_{k} f d\Omega + \int_{\Gamma^{-}_{2}} n_{k} q d\Gamma$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

**Elemental Formulation –** Note that assembly is contained implicitly within the formulation.

$$\sum_{i=1}^{N} \left( A_{ki} u_{i} + B_{ki} \ddot{u}_{i} \right) = F_{k}(t) \qquad k = 1, 2, ..., N$$
$$A_{ki} = \iint_{\Omega^{-}} \left( \frac{\partial n_{k}}{\partial x} k \frac{\partial n_{i}}{\partial x} + \frac{\partial n_{k}}{\partial y} k \frac{\partial n_{i}}{\partial y} \right) d\Omega + \int_{\Gamma^{-}} n_{k} h n_{i} d\Gamma$$
$$B_{ki} = \iint_{\Omega^{-}} n_{k} \rho n_{i} d\Omega$$
$$F_{k} = \iint_{\Omega^{-}} n_{k} f d\Omega + \int_{\Gamma^{-}_{2}} n_{k} q d\Gamma$$

# **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

**Elemental Formulation** – In terms of the corresponding elementally based interpolations

$$u_e(x, y) = \mathbf{N}^T \mathbf{u}_e = \mathbf{u}_e^T \mathbf{N}$$

The finite element model can be expressed as

 $Au + B\ddot{u} = F$ 

$$\mathbf{A} = \sum_{e} \mathbf{k}_{G} + \sum_{e} \mathbf{a}_{G} \qquad \mathbf{B} = \sum_{e} \mathbf{r}_{G}$$
$$\mathbf{F} = \sum_{e} \mathbf{f}_{G} + \sum_{e} \mathbf{q}_{G}$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

Elemental Formulation – In terms of the corresponding elementally based interpolations

$$u_e(x, y) = \mathbf{N}^T \mathbf{u}_e = \mathbf{u}_e^T \mathbf{N}$$

The finite element model can be expressed as

$$\mathbf{k}_{e} = \iint_{A_{e}} \left( \frac{\partial \mathbf{N}}{\partial x} k \frac{\partial \mathbf{N}^{T}}{\partial x} + \frac{\partial \mathbf{N}}{\partial y} k \frac{\partial \mathbf{N}^{T}}{\partial y} \right) dA$$
$$\mathbf{r}_{e} = \iint_{A_{e}} \mathbf{N} \rho \mathbf{N}^{T} dA \qquad \mathbf{a}_{e} = \int_{\gamma_{2e}} \mathbf{N} \alpha \mathbf{N}^{T} ds$$
$$\mathbf{f}_{e} = \iint_{A_{e}} \mathbf{N} f dA \qquad \mathbf{q}_{e} = \int_{\gamma_{2e}} \mathbf{N} q ds$$

#### **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

- **Elemental Formulation** The initial conditions for the system of first-order differential equations are obtained from the initial conditions prescribed for the original initial-boundary value problem.
- Generally  $\mathbf{u}(0)$  is determined by evaluating the function c(x, y) at the nodes to obtain:

 $\mathbf{u}(0) = \mathbf{u}_0 = \begin{pmatrix} \mathbf{c}_1 & \mathbf{c}_2 & \mathbf{c}_3 & \dots & \mathbf{c}_{N-1} & \mathbf{c}_N \end{pmatrix}^{\mathrm{T}}$ 

where  $c_i = c(x_i, y_i)$  with  $(x_i, y_i)$  the coordinates of the *i*<sup>th</sup> node.

#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

- **Elemental Formulation –** The initial conditions for the system of first-order differential equations are obtained from the initial conditions prescribed for the original initial-boundary value problem.
- Generally  $\mathbf{u}(0)$  is determined by evaluating the function d(x, y) at the nodes to obtain

$$\dot{\mathbf{u}}(0) = \dot{\mathbf{u}}_0 = \langle d_1 \quad d_2 \quad d_3 \quad \dots \quad d_{N-1} \quad d_N \rangle^{T}$$

where  $d_i = d(x_i, y_i)$  with  $(x_i, y_i)$  the coordinates of the *i*<sup>th</sup> node.

# **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

**Constraints** - The constraints arise from the boundary conditions specified on  $\Gamma_1$ .

- Generally the values of the constraints are determined from the *q* function with the constrained value of *u* at a node on  $\Gamma_1$  being taken as the value of *q* at that point.
- These constraints are then enforced on the assembled equations, resulting in the final global constrained set of linear first-order differential equations.

$$M\ddot{u} + Ku = f$$
  $u(0) = u_0$   $\dot{u}(0) = \dot{u}_0$ 

# TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

- **Solution** The system of equations is precisely the same in form and character as the corresponding equations developed in Section 4.2.2 for the one-dimensional wave problem.
- The analytical method as well as the numerical methods using the central difference and Newmark algorithms can be used for integrating the above set of equations.

#### **Two-Dimensional Wave Equations**

# The Galerkin Finite Element Method

**Solution** - The central difference algorithm will be conditionally stable with the critical time step depending on the maximum eigenvalue of the associated problem  $(\mathbf{K} - \omega^2 \mathbf{M})\mathbf{v} = 0.$ 

The Newmark algorithm will be unconditionally stable for  $\delta = 0.5$  and  $\alpha = 0.25(\delta + 0.5)^2$ .

# TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations**

## The Galerkin Finite Element Method

**Derived variables** - In a physical situation governed by a wave equation, the derived variables are usually the internal forces computed according to  $F_e = k_e u_e$  per element for each time step.

#### **Two-Dimensional Wave Equations - Example 1**

Consider the two-dimensional wave problem shown below.



#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations - Example 1**

**Discretization.** Due to symmetry, we will model the top right-most quadrant of the membrane using four equally-sized 4-noded quadrilaterals.



# **Two-Dimensional Wave Equations - Example 1**

**Elemental Formulation** – This *N* x *N* set of linear algebraic equations can be written as:

$$\sum_{i=1}^{N} \left( A_{ki} \phi_{i} + \psi B_{ki} \ddot{\phi}_{i} \right) = F_{k}(t) \qquad k = 1, 2, ..., N$$
$$A_{ki} = \iint_{\Omega^{-}} \left( \frac{\partial n_{k}}{\partial x} \frac{\partial n_{i}}{\partial x} + \frac{\partial n_{k}}{\partial y} \frac{\partial n_{i}}{\partial y} \right) d\Omega$$
$$B_{ki} = \iint_{\Omega^{-}} n_{k} \psi n_{i} d\Omega$$
$$F_{k} = \iint_{\Omega^{-}} n_{k} f d\Omega$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

**Elemental Formulation** – In terms of the corresponding elementally based interpolations

$$\phi_e(\mathbf{X},\mathbf{Y}) = \mathbf{N}^T \phi_e = \phi_e^T \mathbf{N}$$

The finite element model can be expressed as

$$\mathbf{k}_{e} = \iint_{A_{e}} \left( \frac{\partial \mathbf{N}}{\partial x} \frac{\partial \mathbf{N}^{T}}{\partial x} + \frac{\partial \mathbf{N}}{\partial y} \frac{\partial \mathbf{N}^{T}}{\partial y} \right) dA$$
$$\mathbf{r}_{e} = \iint_{A_{e}} \mathbf{N} \mathbf{N}^{T} dA$$
$$\mathbf{f}_{e} = \iint_{A_{e}} \mathbf{N} f dA$$

# **Two-Dimensional Wave Equations - Example 1**

**Interpolation** - In matrix notation, the distribution of the function over the element in local (*s*, *t*) coordinates is:



# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

**Rectangular Elements -** The elemental interpolations for a rectangular element are:



#### **Two-Dimensional Wave Equations - Example 1**

**Rectangular Elements** - The shape functions are visually deceiving. There is no curvature in directions parallel to any side; however, there is a twist due to the *xy* term in the element representation.



#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations - Example 1**

**Interpolation** - This integrals may be transformed into the local coordinate space using  $as = x - x_0$  and  $bt = y - y_0$ , for a rectangular element  $2a \times 2b$  in size:

$$\iint_{A_{e}} F(x,y) dx dy = \int_{-1}^{1} \int_{-1}^{1} F(x_{0} + as, y_{0} + bt) \frac{\partial(\xi,\eta)}{\partial(s,t)} ds dt$$

where the Jacobian  $\partial(\xi, \eta)/\partial(s, t)$  has a value of  $A_e/4$ .

The resulting global system stiffness matrix  $\mathbf{k}_{e}$  is:

$$\mathbf{k}_{\mathbf{e}} = \frac{1}{6} \begin{bmatrix} 2 & -2 & -1 & 1 \\ -2 & 2 & 1 & -1 \\ -1 & 1 & 2 & -2 \\ 1 & -1 & -2 & 2 \end{bmatrix} + \frac{a}{b} \begin{bmatrix} 2 & 1 & -1 & -2 \\ 1 & 2 & -2 & -1 \\ -1 & -2 & 2 & 1 \\ -2 & -1 & 1 & 2 \end{bmatrix}$$

# **Two-Dimensional Wave Equations - Example 1**

# Rectangular Elements - Evaluation of ke

If the element is square, then a = b, then  $\mathbf{k}_{e}$  becomes:

$$\mathbf{k}_{e} = \frac{1}{6} \begin{bmatrix} \frac{b}{a} \begin{bmatrix} 2 & -2 & -1 & 1\\ -2 & 2 & 1 & -1\\ -1 & 1 & 2 & -2\\ 1 & -1 & -2 & 2 \end{bmatrix} + \frac{a}{b} \begin{bmatrix} 2 & 1 & -1 & -2\\ 1 & 2 & -2 & -1\\ -1 & -2 & 2 & 1\\ -2 & -1 & 1 & 2 \end{bmatrix}$$
$$\mathbf{k}_{e} = \frac{1}{6} \begin{bmatrix} 4 & -1 & -2 & -1\\ -1 & 4 & -1 & -2\\ -2 & -1 & 4 & -1\\ -1 & -2 & -1 & 4 \end{bmatrix}$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

The elemental  $\mathbf{m}_{e}$  matrices are:

$$\mathbf{m}_{\mathbf{e}} = \iint_{A_{\mathbf{e}}} \mathbf{N} \mathbf{N}^{\mathsf{T}} \, dA = \int_{-1}^{1} \int_{-1}^{1} F(\mathbf{x}_{0} + \mathbf{a}\mathbf{s}, \mathbf{y}_{0} + \mathbf{b}t) \frac{\partial(\xi, \eta)}{\partial(\mathbf{s}, t)} d\mathbf{s} \, dt$$

where the Jacobian  $\partial(\xi, \eta)/\partial(s, t)$  has a value of  $A_e/4$ .

The resulting global system stiffness matrix  $\mathbf{m}_{\mathbf{e}}$  is:

$$\mathbf{m}_{e} = \frac{A_{e}}{36} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 4 & 2 & 1 \\ 1 & 2 & 4 & 2 \\ 2 & 1 & 2 & 4 \end{bmatrix}$$

#### **Two-Dimensional Wave Equations - Example 1**

Transforming the integral into the non-dimensional coordinates (s, t) yields:

$$\mathbf{f_{e}} \approx \left( \int_{-1}^{1} \int_{-1}^{1} \mathbf{NN^{T}} \frac{A_{e}}{4} ds dt \right) \mathbf{f}$$
$$\mathbf{f_{e}} \approx \frac{A_{e}}{36} \begin{cases} 4f_{1} + 2f_{2} + f_{3} + 2f_{4} \\ 2f_{1} + 4f_{2} + 2f_{3} + f_{4} \\ f_{1} + 2f_{2} + 4f_{3} + 2f_{4} \\ 2f_{1} + f_{2} + 2f_{3} + 4f_{4} \end{cases}$$

The resulting 4 x 1 elemental load vector contributes to the global system equations at those locations corresponding to the four nodes defining the element.

# TIME-DEPENDENT PROBLEMS





# TIME-DEPENDENT PROBLEMS





#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations - Example 1**

The assembled  $\mathbf{K}_{\mathbf{G}}$ ,  $\mathbf{M}_{\mathbf{G}}$ , and  $\mathbf{F}$  matrices are:

$$\mathbf{K}_{\mathbf{G}} = \frac{1}{6} \begin{bmatrix} 4 & -1 & -2 & -1 & 0 & 0 & 0 & 0 & 0 \\ -1 & 8 & -2 & -2 & -1 & -2 & 0 & 0 & 0 \\ -2 & -2 & 16 & -2 & -2 & -2 & -2 & -2 & -2 \\ -1 & -2 & -2 & 8 & 0 & 0 & 0 & -2 & -1 \\ 0 & -1 & -2 & 0 & 4 & -1 & 0 & 0 & 0 \\ 0 & -2 & -2 & 0 & -1 & 8 & -1 & -2 & 0 \\ 0 & 0 & -2 & -2 & 0 & -1 & 4 & -1 & 0 \\ 0 & 0 & -2 & -2 & 0 & -2 & -1 & 8 & -1 \\ 0 & 0 & -2 & -1 & 0 & 0 & 0 & -1 & 4 \end{bmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \\ \phi_5 \\ \phi_6 \\ \phi_7 \\ \phi_8 \\ \phi_9 \end{bmatrix}$$

# **Two-Dimensional Wave Equations - Example 1**

The assembled  $\mathbf{K}_{\mathbf{G}}$ ,  $\mathbf{M}_{\mathbf{G}}$ , and  $\mathbf{F}$  matrices are:

$$\mathbf{M}_{\mathbf{G}} = \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 & 0 & 0 & 0 & 0 & 0 \\ 2 & 8 & 4 & 1 & 2 & 1 & 0 & 0 & 0 \\ 1 & 4 & 16 & 4 & 1 & 4 & 1 & 4 & 1 \\ 2 & 1 & 4 & 8 & 0 & 0 & 0 & 1 & 2 \\ 0 & 2 & 1 & 0 & 4 & 2 & 0 & 0 & 0 \\ 0 & 1 & 4 & 0 & 2 & 8 & 2 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 & 2 & 4 & 2 & 0 \\ 0 & 0 & 4 & 1 & 0 & 1 & 2 & 8 & 2 \\ 0 & 0 & 1 & 2 & 0 & 0 & 0 & 2 & 4 \end{bmatrix} \begin{array}{c} \phi_{1} \\ \phi_{2} \\ \phi_{3} \\ \phi_{4} \\ \phi_{5} \\ \phi_{6} \\ \phi_{7} \\ 0.01880 \\ 0.00347 \\ 0.01880 \\ 0.01329 \end{bmatrix}$$

## TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

Applying nonhomogeneous boundary conditions that  $\{\phi_5, \dots, \phi_9\} = \{0, 0, 0, 0, 0\}$ ; the constrained equations are:

$$\mathbf{K}_{\mathbf{G}}\boldsymbol{\phi} + \boldsymbol{\psi}\mathbf{M}_{\mathbf{G}}\boldsymbol{\ddot{\phi}} = \mathbf{F}$$

|   | 4  | -1 | -2 | -1 | 0  | 0  | 0  | 0  | 0  | $\left[\phi_{1}\right]$ | )  | 4 | 2  | 1 | 2 | 0  | 0 | 0 | 0 | 0                      | $\left( \ddot{\phi}_{1} \right)$ |                  | 0.05089         | )       |         |   |
|---|----|----|----|----|----|----|----|----|----|-------------------------|--|---|--|---|---|----|---|---|---|------------------------|----------------------------------|------------------|-----------------|---------|---------|---|
|   | -1 | 8  | -2 | -2 | -1 | -2 | 0  | 0  | 0  | $\phi_2$                | $\phi_2$ $\phi_3$ $\phi_4$ $\phi_4$ $\phi_4$ |   | 2  | 8 | 4 | 1  | 2 | 1 | 0 | 0                      | 0                                | $\ddot{\phi}_2$  |                 | 0.07197 | ł       |   |
|   | -2 | -2 | 16 | -2 | -2 | -2 | -2 | -2 | -2 | $\phi_3$                |  |   | $\left  \begin{array}{c} \phi_3 \\ \phi_4 \end{array} \right  = 2$ | 1 | 4 | 16 | 4 | 1 | 4 | 1                      | 4                                | 1                | $\ddot{\phi}_3$ |         | 0.10178 | l |
| 4 | -1 | -2 | -2 | 8  | 0  | 0  | 0  | -2 | -1 | $\phi_4$                |  |   |  | 2 | 1 | 4  | 8 | 0 | 0 | 0                      | 1                                | 2                | $\ddot{\phi}_4$ |         | 0.07197 | l |
| 1 | 0  | -1 | -2 | 0  | 4  | -1 | 0  | 0  | 0  | 6                       | $\left\{+\frac{\psi}{144}\right\}$           | 0 | 2  | 1 | 0 | 4  | 2 | 0 | 0 | 0                      | ¢₅                               | } = <            | 0.01329         | ł       |         |   |
| ь | 0  | -2 | -2 | 0  | -1 | 8  | -1 | -2 | 0  | 0                       | 144  | 0 | 1  | 4 | 0 | 2  | 8 | 2 | 1 | 0                      | $\dot{\phi}_6$                   |                  | 0.01880         | ł       |         |   |
|   | 0  | 0  | -2 | 0  | 0  | -1 | 4  | -1 | 0  | 0                       |  |   | 0  | 0 | 1 | 0  | 0 | 2 | 4 | 2                      | 0                                | ø,               |                 | 0.00347 | l       |   |
|   | 0  | 0  | -2 | -2 | 0  | -2 | -1 | 8  | -1 | 0                       |  |   | 0  | 0 | 4 | 1  | 0 | 1 | 2 | 8                      | 2                                | $\dot{\phi}_{B}$ |                 | 0.01880 | ł       |   |
|   | 0  | 0  | -2 | -1 | 0  | 0  | 0  | -1 | 4  | 0                       | J  | 0 | 0  | 1 | 2 | 0  | 0 | 0 | 2 | 4                      | ,                                |                  | 0.01329         | J       |         |   |
|   |    |    |    |    |    |    |    |    |    |                         |  |   |  |   |   |    |   |   |   |                        | J                                |                  |                 |         |         |   |
|   |    |    |    |    |    |    |    |    |    |                         |  |   |  |   |   |    |   |   | { | $\ddot{\phi}_5, \cdot$ | ,¢                               | ;<br>;}=         | = 0             |         |         |   |
|   |    |    |    |    |    |    |    |    |    |                         |  |   |  |   |   |    |   |   |   |                        |                                  |                  |                 |         |         |   |

# **Two-Dimensional Wave Equations - Example 1**

**Assembly.** With both the boundary conditions essential, **BT**=0 and **bt**=0. It follows that the assembled equations are:

$$\mathbf{K}_{\mathbf{G}}\boldsymbol{\phi} + \boldsymbol{\psi}\mathbf{M}_{\mathbf{G}}\boldsymbol{\ddot{\phi}} = \mathbf{F}$$

|   | 4  | -1 | -2 | -1] | $\left( \phi_{1} \right)$ |     | 4 | 2 | 1  | 2] | $\left( \ddot{\phi}_{1} \right)$ |      | 0.05089 |
|---|----|----|----|-----|---------------------------|-----|---|---|----|----|----------------------------------|------|---------|
| 1 | -1 | 8  | -2 | -2  | $\phi_2$                  | ļψ  | 2 | 8 | 4  | 1  | $ \ddot{\phi}_2 $                |      | 0.07197 |
| 6 | -2 | -2 | 16 | -2  | $\phi_3$                  | 144 | 1 | 4 | 16 | 4  | $\vec{\phi}_3$                   | >= < | 0.10178 |
|   | 1  | -2 | -2 | 8   | $\phi_4$                  |     | 2 | 1 | 4  | 8  | $\left  \ddot{\phi}_{4} \right $ |      | 0.07197 |

Let assume the  $\psi$  = 1 for this example.

# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

For the Newmark algorithm, the critical time step is associated with the largest eigenvalue of the  $(\mathbf{K} - \omega^2 \mathbf{M})\phi = 0$  system.

$$h_{cr} < \frac{2}{\omega_{max}}$$

The eigenvalue problem is:

$$\begin{pmatrix} \mathbf{K}_{\mathbf{G}} + \omega^{2} \mathbf{M}_{\mathbf{G}} \end{pmatrix} \phi = \mathbf{0}$$

$$\begin{pmatrix} \frac{1}{6} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 8 & -2 & -2 \\ -2 & -2 & 16 & -2 \\ -1 & -2 & -2 & 8 \end{bmatrix} - \frac{\omega^{2}}{144} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 8 & 4 & 1 \\ 1 & 4 & 16 & 4 \\ 2 & 1 & 4 & 8 \end{bmatrix} \phi = \mathbf{0}$$

#### 24/47

# TIME-DEPENDENT PROBLEMS

## **Two-Dimensional Wave Equations - Example 1**

The eigenvalues can be found using a variety of available solution techniques. In Matlab, use [V,D] = eig(A).

For example,  $[V,D] = eig(M_G \setminus K_G)$  gives eigenvalues D:

 $\omega_1^2 = 5.19332$   $\omega_2^2 = 34.28571$   $\omega_3^2 = 34.28571$   $\omega_4^2 = 63.37811$ 

The eigenvectors v are:

|                           | 0.66667 | 1                        | 0.89443   |                               | 0.09022  |  | –0.66667) |
|---------------------------|---------|--------------------------|-----------|-------------------------------|----------|--|-----------|
| V -                       | 0.47140 | V _                      | 0.00000   |                               | -0.70350 | V -                                    | 0.47140   |
| <b>v</b> <sub>1</sub> = - | 0.33333 | $\mathbf{v}_2 = \langle$ | ]–0.44721 | $\rightarrow$ $V_3 = \langle$ | -0.04511 | $\rightarrow$ $\mathbf{v}_4 = \langle$ | -0.33333  |
|                           | 0.47140 |                          | 0.00000   |                               | 0.70350  |  | 0.47140   |

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

For the example the largest eigenvalue is 63.37811, the critical time step is:

$$h_{cr} < \frac{2}{\omega_{max}} < \frac{2}{\sqrt{63.37811}} < 0.2512 \, \text{sec}$$

# **Two-Dimensional Wave Equations - Example 1**

**Newmark's Method**– The equations to be solved at the first step can be written as:

$$\left(\mathbf{M} + \alpha h^{2}\mathbf{K}\right)\phi_{1} = \mathbf{M}\left(\phi_{0} + h\dot{\phi}_{0} + c_{1}h^{2}\ddot{\phi}_{0}\right) + \alpha h^{2}\mathbf{F}_{1}$$

$$\dot{\phi}_{1} = \dot{\phi}_{0} + \frac{\delta\left(\phi_{1} - \phi_{0} - h\dot{\phi}_{0}\right)}{\alpha h} + c_{2}h\ddot{\phi}_{0}$$

$$\ddot{\phi}_{1} = \frac{\phi_{1} - \phi_{0} - h\dot{\phi}_{0}}{\alpha h^{2}} - \frac{c_{1}\ddot{\phi}_{0}}{\alpha}$$

$$c_{1} = \frac{1}{2} - \alpha$$

$$c_{2} = 1 - \frac{\delta}{2\alpha}$$

# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

**Newmark's Method**– Evaluating the differential equation at t = 0 yields

$$\mathbf{M}\ddot{\phi}_{0} = \mathbf{F} - \mathbf{K}\phi_{0} \implies \ddot{\phi}_{0} = \mathbf{M}^{-1} \left(\mathbf{F} - \mathbf{K}\phi_{0}\right)^{(1)} \begin{bmatrix} 0.05089 \\ 0.07197 \\ 0.07711 \\ 0.50000 \\ 0.70711 \\ 0.50000 \\ 0.70711 \\ 0.50000 \\ 0.70711 \\ 0.50000 \\ 0.70711 \\ 0.0000 \\ 0.$$

#### **Two-Dimensional Wave Equations - Example 1**

**Newmark's Method**– Taking  $\alpha = 0.25$ ,  $\delta = 0.5$ , and h = 0.01 seconds yields:



#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations - Example 1**

**Newmark's Method**– The solution for  $\phi_1$  is

 $\phi_1 = 10^{-5} \langle 4.9994$  3.5351 2.4997 3.5351 $\rangle^{T}$ 

Now solve for the velocity at t = h

$$\dot{\phi}_{1} = \phi_{0}^{\dagger} + \frac{\delta\left(\phi_{1} - \phi_{0} - h\dot{\phi}_{0}\right)^{\dagger}}{\alpha h} + c_{2}h\ddot{\phi}_{0}^{\dagger}$$
$$\dot{\phi}_{1} = \frac{\delta\phi_{1}}{\alpha h} + c_{2}h\ddot{\phi}_{0} = \frac{\delta}{\alpha h} \begin{cases} 4.99935\\ 3.53507\\ 2.49968\\ 3.53507 \end{cases} 10^{-5} + c_{2}h \begin{cases} 1.00000\\ 0.70711\\ 0.50000\\ 0.70711 \end{cases}$$
$$= \langle 0.01000 \quad 0.00707 \quad 0.005 \quad 0.00707 \rangle^{T}$$

#### 27/47

# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

Newmark's Method- The equation for acceleration at

$$t = 0 \text{ is } \phi_0(0) = 0$$
  
$$\ddot{\phi}_1 = \frac{\phi_1 - \phi_0 - h\dot{\phi}_0}{\alpha h^2} - \frac{c_1 \ddot{\phi}_0}{\alpha} = \frac{\phi_1}{\alpha h^2} - \frac{c_1 \ddot{\phi}_0}{\alpha}$$

Now solve for the acceleration at t = h

|                                  | (4.99935) |                      | [1.00000] |  |
|----------------------------------|-----------|----------------------|-----------|--|
| $\frac{1}{2}\delta$              | 3.53507   | 10 <sup>-5</sup> a b | 0.70711   |  |
| $\varphi_1 = \frac{1}{\alpha h}$ | 2.49968   | $>10 - C_2 II <$     | 0.50000   |  |
|                                  | 3.53507   |                      | 0.70711   |  |
|                                  |           |                      |           |  |

$$\ddot{\phi}_1 = \left< 0.99974 \quad 0.70692 \quad 0.49987 \quad 0.70692 \right>^7$$

#### TIME-DEPENDENT PROBLEMS

Two-Dimensional Wave Equations - Example 1 Newmark's Method – For n = 2:

$$\phi_2 = 10^{-4} \langle 1.99948 \quad 1.41385 \quad 0.99974 \quad 1.41385 \rangle^7$$
  
 $\dot{\phi}_2 = \langle 0.01999 \quad 0.01414 \quad 0.01000 \quad 0.01414 \rangle^T$   
 $\ddot{\phi}_2 = \langle 0.99896 \quad 0.70637 \quad 0.49948 \quad 0.70637 \rangle^T$ 

For n = 3:

 $\phi_3 = 10^{-4} \langle 4.49786 \quad 3.18047 \quad 2.24893 \quad 3.18047 \rangle^{\tau}$  $\dot{\phi}_3 = \langle 0.02998 \quad 0.02120 \quad 0.01499 \quad 0.02120 \rangle^{\tau}$  $\ddot{\phi}_3 = \langle 0.99766 \quad 0.70546 \quad 0.49883 \quad 0.70546 \rangle^{\tau}$ 

# **Two-Dimensional Wave Equations - Example 1**

The values of  $\phi_1$ ,  $\phi_2$ , and  $\phi_3$  for 0 < t < 10 sec. are show below:



## TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

The velocity values of  $\phi_1$ ,  $\phi_2$ , and  $\phi_3$  for  $0 \le t \le 10$  sec. are:



# **Two-Dimensional Wave Equations - Example 1**

The acceleration values of  $\phi_1$ ,  $\phi_2$ , and  $\phi_3$  for 0 < t < 10 sec. are:



# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**



# **Two-Dimensional Wave Equations - Example 1**



# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

Consider the two-dimensional wave problem shown below.



## **Two-Dimensional Wave Equations - Example 2**

**Discretization.** Due to symmetry, we will model the top right-most quadrant of the membrane using four equally-sized 4-noded quadrilaterals.



#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

# Rectangular Elements - Evaluation of k<sub>e</sub>

As in Example 1, all the  $\mathbf{k}_{e}$  are:

$$\mathbf{k}_{1} = \frac{1}{6} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ 2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 4 & -1 & -2 \\ -2 & -1 & 4 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & -2 & -1 & 4 \end{bmatrix} \end{bmatrix} \end{bmatrix} \begin{bmatrix}$$

#### 32/47

## TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

The elemental  $\mathbf{m}_{\mathbf{e}}$  matrices are:

$$\mathbf{m}_{\mathbf{e}} = \iint_{A_{\mathbf{e}}} \mathbf{N} \mathbf{N}^{\mathsf{T}} \, dA = \int_{-1}^{1} \int_{-1}^{1} F(\mathbf{x}_{0} + \mathbf{a}\mathbf{s}, \mathbf{y}_{0} + \mathbf{b}t) \frac{\partial(\xi, \eta)}{\partial(\mathbf{s}, t)} d\mathbf{s} \, dt$$

where the Jacobian  $\partial(\xi, \eta)/\partial(s, t)$  has a value of  $A_e/4$ .

The resulting global system stiffness matrix  $\mathbf{m}_{\mathbf{e}}$  is:

 $\mathbf{m}_{\mathbf{e}} = \frac{A_{\mathbf{e}}}{36} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 4 & 2 & 1 \\ 1 & 2 & 4 & 2 \\ 2 & 1 & 2 & 4 \end{bmatrix}$ 

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

Rectangular Elements - Evaluation of m<sub>e</sub>

As in Example 1, all the  $m_e$  are:

$$\mathbf{m}_{1} = \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 4 & 2 & 1 \\ 1 & 2 & 4 & 2 \\ 2 & 1 & 2 & 4 \end{bmatrix} \begin{bmatrix} 1 \\ 2 \\ 3 \\ 4 \end{bmatrix} \qquad \mathbf{m}_{2} = \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 4 & 2 & 1 \\ 1 & 2 & 4 & 2 \\ 2 & 1 & 2 & 4 \end{bmatrix} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 5 \\ 6 \\ 2 \\ 1 & 2 & 4 \end{bmatrix} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 \\ 1 & 2 & 4 \end{bmatrix} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 3 \\ 3 \\ 3 \end{bmatrix}$$
$$\mathbf{m}_{3} = \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 4 & 2 & 1 \\ 1 & 2 & 4 & 2 \\ 2 & 1 & 2 & 4 \end{bmatrix} \begin{bmatrix} 3 \\ 6 \\ 7 \\ 8 \end{bmatrix} \qquad \mathbf{m}_{4} = \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 4 & 2 & 1 \\ 1 & 2 & 4 & 2 \\ 2 & 1 & 2 & 4 \end{bmatrix} \begin{bmatrix} 3 \\ 8 \\ 9 \\ 4 \end{bmatrix}$$

# **Two-Dimensional Wave Equations - Example 2**

Transforming the integral into the non-dimensional coordinates (s, t) yields:

$$\mathbf{f}_{\mathbf{e}} \approx \left(\int_{-1}^{1} \int_{-1}^{1} \mathbf{N} \mathbf{N}^{\mathsf{T}} \frac{A_{e}}{4} ds dt\right) \mathbf{f}$$
$$\mathbf{f}_{\mathbf{e}} \approx \frac{A_{e}}{36} \begin{cases} 4f_{1} + 2f_{2} + f_{3} + 2f_{4} \\ 2f_{1} + 4f_{2} + 2f_{3} + f_{4} \\ f_{1} + 2f_{2} + 4f_{3} + 2f_{4} \\ 2f_{1} + f_{2} + 2f_{3} + 4f_{4} \end{cases}$$

The resulting 4 x 1 elemental load vector contributes to the global system equations at those locations corresponding to the four nodes defining the element.

# TIME-DEPENDENT PROBLEMS





# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

The assembled  $\mathbf{K}_{\mathbf{G}}$ ,  $\mathbf{M}_{\mathbf{G}}$ , and  $\mathbf{F}$  matrices are:

$$\mathbf{K}_{\mathbf{G}} = \frac{1}{6} \begin{bmatrix} 4 & -1 & -2 & -1 & 0 & 0 & 0 & 0 & 0 \\ -1 & 8 & -2 & -2 & -1 & -2 & 0 & 0 & 0 \\ -2 & -2 & 16 & -2 & -2 & -2 & -2 & -2 & -2 \\ -1 & -2 & -2 & 8 & 0 & 0 & 0 & -2 & -1 \\ 0 & -1 & -2 & 0 & 4 & -1 & 0 & 0 & 0 \\ 0 & -2 & -2 & 0 & -1 & 8 & -1 & -2 & 0 \\ 0 & 0 & -2 & 0 & 0 & -1 & 4 & -1 & 0 \\ 0 & 0 & -2 & -2 & 0 & -2 & -1 & 8 & -1 \\ 0 & 0 & -2 & -1 & 0 & 0 & 0 & -1 & 4 \end{bmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \\ \phi_5 \\ \phi_6 \\ \phi_7 \\ \phi_8 \\ \phi_9 \end{bmatrix}$$

# **Two-Dimensional Wave Equations - Example 2**

The assembled  $\mathbf{K}_{\mathbf{G}},\,\mathbf{M}_{\mathbf{G}},\,\text{and}\,\,\mathbf{F}$  matrices are:

$$\mathbf{M}_{\mathbf{G}} = \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 & 0 & 0 & 0 & 0 & 0 \\ 2 & 8 & 4 & 1 & 2 & 1 & 0 & 0 & 0 \\ 1 & 4 & 16 & 4 & 1 & 4 & 1 & 4 & 1 \\ 2 & 1 & 4 & 8 & 0 & 0 & 0 & 1 & 2 \\ 0 & 2 & 1 & 0 & 4 & 2 & 0 & 0 & 0 \\ 0 & 1 & 4 & 0 & 2 & 8 & 2 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 & 2 & 4 & 2 & 0 \\ 0 & 0 & 4 & 1 & 0 & 1 & 2 & 8 & 2 \\ 0 & 0 & 1 & 2 & 0 & 0 & 0 & 2 & 4 \end{bmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_6 \\ \phi_7 \\ \phi_8 \\ \phi_9 \end{bmatrix}$$

## TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

The assembled  $\mathbf{K}_{\mathbf{G}},\,\mathbf{M}_{\mathbf{G}},\,\text{and}\;\mathbf{F}$  matrices are:

|              | 0.02778 | $\phi_1$         |              | (0) | $\phi_1$                                |
|--------------|---------|------------------|--------------|-----|---|
|              | 0.01389 | $\phi_2$         |              | 0   | $\phi_2$                                |
|              | 0.00694 | $\phi_3$         |              | 0   | $\phi_3$                                |
|              | 0.01389 | $\phi_4$         |              | 0   | $\phi_4$                                |
| <b>F</b> = < | 0.00000 | ¢5               | <b>F</b> = < | 0   | · • • • • • • • • • • • • • • • • • • • |
|              | 0.00000 | $\phi_6$         |              | 0   | $\phi_6$                                |
|              | 0.00000 | $\phi_7$         |              | 0   | $\phi_7$                                |
|              | 0.00000 | $\phi_8$         |              | 0   | $\phi_8$                                |
|              | 0.00000 | t≤0.1 <b>¢</b> 9 |              | 0   | <sub>t&gt;0.1</sub>                     |

# **Two-Dimensional Wave Equations - Example 2**

Applying nonhomogeneous boundary conditions that  $\{\phi_5, \dots, \phi_9\} = \{0, 0, 0, 0, 0\}$ ; the constrained equations are:

$$\mathbf{K}_{\mathbf{G}}\phi + \mathbf{M}_{\mathbf{G}}\ddot{\phi} = \mathbf{F}$$

|    |  |  |   |  |   |  |  |  |  |   | _   |   |   |   |   |   |   |   | _   | <i></i>   |   |   |   |         |
|----|--|--|---|--|---|--|--|--|--|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---------|
| 4  | -1   | -2   | -1  | 0  | 0   | 0  | 0  | 0  | $\phi_1$   |   | 4   | 2   | 1   | 2   | 0   | 0   | 0   | 0   | 0   | $\phi_1$  |   | 0.02778   |   |         |
| -1 | 8  | -2   | -2  | -1   | -2  | 0  | 0  | 0  | $\phi_2$   |   | 2   | 8   | 4   | 1   | 2   | 1   | 0   | 0   | 0   | $\ddot{\phi}_2$                                       |   | 0.01389   |   |         |
| -2 | -2   | 16   | -2  | -2   | -2  | -2   | -2   | -2   | $\phi_3$   |   | 1   | 1   | 4   | 16  | 4   | 1   | 4   | 1   | 4   | 1   | $\ddot{\phi}_3$                                       |   | 0.00694   |         |
| -1 | -2   | -2   | 8   | 0  | 0   | 0  | -2   | -1   | $\phi_4$   |   |   | 2   | 1   | 4   | 8   | 0   | 0   | 0   | 1   | 2   | $\ddot{\phi}_4$                                       |   | 0.01389   |         |
| 0  | -1   | -2   | 0   | 4  | -1  | 0  | 0  | 0  | 0  | $+\frac{1}{144}$                                      | 0   | 2   | 1   | 0   | 4   | 2   | 0   | 0   | 0   | { <b>∅</b> ₅  | } = <   | 0.00000   |   |         |
| 0  | -2   | -2   | 0   | -1   | 8   | -1   | -2   | 0  | 0  | 144   | 0   | 1   | 4   | 0   | 2   | 8   | 2   | 1   | 0   | $\dot{\phi}_{6}$                                      |   | 0.00000   |   |         |
| 0  | 0  | -2   | 0   | 0  | -1  | 4  | -1   | 0  | 0  |   |   | 0   | 0   | 1   | 0   | 0   | 2   | 4   | 2   | 0   | ø,  |   | 0.00000   |         |
| 0  | 0  | -2   | -2  | 0  | -2  | -1   | 8  | -1   | 0  |   |   |   | 0   | 0   | 4   | 1   | 0   | 1   | 2   | 8   | 2   | ∮ <sub>8</sub>  |   | 0.00000 |
| 0  | 0  | -2   | -1  | 0  | 0   | 0  | -1   | 4  | 0  |   | 0   | 0   | 1   | 2   | 0   | 0   | 0   | 2   | 4   | ø,  | J   | 0.00000   | <i>t</i> ≤0.1   |         |
|    |  |  |   |  |   |  |  |  |  |   |   |   |   |   |   |   |   |   |   | ♦   |   |   |   |         |
|    |  |  |   |  |   |  |  |  |  |   |   |   |   |   |   |   |   | {   | $\ddot{\phi}_5, \cdot$                                | , ģ   | ₿<br> <br> <br> <br> <br>                             | = 0   |   |         |
|    | 4<br>-1<br>-2<br>-1<br>0<br>0<br>0<br>0<br>0 | 4 -1<br>-1 8<br>-2 -2<br>-1 -2<br>0 -1<br>0 -2<br>0 0<br>0 0<br>0 0<br>0 0 | 4     -1     -2       -1     8     -2       -2     -2     16       -1     -2     -2       0     -1     -2       0     -2     -2       0     0     -2       0     0     -2       0     0     -2       0     0     -2 | 4       -1       -2       -1         -1       8       -2       -2         -2       -2       16       -2         -1       -2       -2       8         0       -1       -2       0         0       -2       -2       0         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -2         0       0       -2       -1 | 4     -1     -2     -1     0       -1     8     -2     -2     -1       -2     -2     16     -2     -2       -1     -2     -2     8     0       0     -1     -2     0     4       0     -2     -2     0     -1       0     0     -2     -2     0     0       0     0     -2     -2     0     0       0     0     -2     -2     0     0       0     0     -2     -2     0     0       0     0     -2     -1     0 | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | 4     -1     -2     -1     0     0       -1     8     -2     -2     -1     -2     0       -2     -2     16     -2     -2     -2     -2       -1     -2     2     8     0     0     0       0     -1     -2     0     4     -1     0       0     -2     -2     0     -1     8     -1       0     0     -2     0     -1     4       0     0     -2     -2     0     -1       0     0     -2     -2     0     -2     -1 | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $\begin{array}{cccccccccccccccccccccccccccccccccccc$ | $ \begin{array}{cccccccccccccccccccccccccccccccccccc$ |         |

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

**Assembly.** With both the boundary conditions essential, **BT**=0 and **bt**=0. It follows that the assembled equations are:

$$\mathbf{K}_{\mathbf{G}}\phi + \mathbf{M}_{\mathbf{G}}\ddot{\phi} = \mathbf{F}$$

$$\frac{1}{6}\begin{bmatrix}4 & -1 & -2 & -1\\-1 & 8 & -2 & -2\\-2 & -2 & 16 & -2\\-1 & -2 & -2 & 8\end{bmatrix}\begin{pmatrix}\phi_{1}\\\phi_{2}\\\phi_{3}\\\phi_{4}\end{bmatrix} + \frac{1}{144}\begin{bmatrix}4 & 2 & 1 & 2\\2 & 8 & 4 & 1\\1 & 4 & 16 & 4\\2 & 1 & 4 & 8\end{bmatrix}\begin{bmatrix}\ddot{\phi}_{1}\\\ddot{\phi}_{2}\\\ddot{\phi}_{3}\\\ddot{\phi}_{4}\end{bmatrix} = \begin{bmatrix}0.027778\\0.013889\\0.006944\\0.013889\end{bmatrix}_{t\leq 0.1}$$
$$\frac{1}{6}\begin{bmatrix}4 & -1 & -2 & -1\\-1 & 8 & -2 & -2\\-2 & -2 & 16 & -2\\-1 & -2 & -2 & 8\end{bmatrix}\begin{bmatrix}\phi_{1}\\\phi_{2}\\\phi_{3}\\\phi_{4}\end{bmatrix} + \frac{1}{144}\begin{bmatrix}4 & 2 & 1 & 2\\2 & 8 & 4 & 1\\1 & 4 & 16 & 4\\2 & 1 & 4 & 8\end{bmatrix}\begin{bmatrix}\ddot{\phi}_{1}\\\ddot{\phi}_{2}\\\ddot{\phi}_{3}\\\ddot{\phi}_{4}\end{bmatrix} = \begin{bmatrix}0\\0\\0\\0\\0\\t>0.1$$

## **Two-Dimensional Wave Equations - Example 2**

For the Newmark algorithm, the critical time step is associated with the largest eigenvalue of the  $(\mathbf{K} - \omega^2 \mathbf{M})\phi = 0$  system.

$$h_{cr} < \frac{2}{\omega_{max}}$$

The eigenvalue problem is the same as Example 1:

$$\begin{pmatrix} \mathbf{K}_{\mathbf{G}} + \omega^{2} \mathbf{M}_{\mathbf{G}} \end{pmatrix} \phi = \mathbf{0}$$

$$\begin{pmatrix} \frac{1}{6} \begin{bmatrix} 4 & -1 & -2 & -1 \\ -1 & 8 & -2 & -2 \\ -2 & -2 & 16 & -2 \\ -1 & -2 & -2 & 8 \end{bmatrix} - \frac{1}{144} \begin{bmatrix} 4 & 2 & 1 & 2 \\ 2 & 8 & 4 & 1 \\ 1 & 4 & 16 & 4 \\ 2 & 1 & 4 & 8 \end{bmatrix} \phi = \mathbf{0}$$

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

The eigenvalues can be found using a variety of available solution techniques. In Matlab, use [V,D] = eig(A).

For example,  $[V,D] = eig(M_g \setminus K_g)$  gives eigenvalues D:

$$\omega_1^2 = 5.19332$$
  $\omega_2^2 = 34.28571$   $\omega_3^2 = 34.28571$   $\omega_4^2 = 63.37811$ 

The eigenvectors v are:

$$V_{1} = \begin{cases} 0.66667\\ 0.47140\\ 0.33333\\ 0.47140 \end{cases} \quad V_{2} = \begin{cases} 0.89443\\ 0.00000\\ -0.44721\\ 0.00000 \end{cases} \quad V_{3} = \begin{cases} 0.09022\\ -0.70350\\ -0.04511\\ 0.70350 \end{cases} \quad V_{4} = \begin{cases} -0.66667\\ 0.47140\\ -0.33333\\ 0.47140 \end{cases}$$

# **Two-Dimensional Wave Equations - Example 2**

For the example the largest eigenvalue is 63.37811, the critical time step is:

$$h_{cr} < \frac{2}{\omega_{max}} < \frac{2}{\sqrt{63.37811}} < 0.2512 \, \text{sec}$$

## TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

**Newmark's Method** – The equations to be solved at the first step can be written as:

$$\left(\mathbf{M} + \alpha h^{2}\mathbf{K}\right)\phi_{1} = \mathbf{M}\left(\phi_{0} + h\dot{\phi}_{0} + c_{1}h^{2}\ddot{\phi}_{0}\right) + \alpha h^{2}\mathbf{F}_{1}$$

$$\dot{\phi}_{1} = \dot{\phi}_{0} + \frac{\delta\left(\phi_{1} - \phi_{0} - h\dot{\phi}_{0}\right)}{\alpha h} + c_{2}h\ddot{\phi}_{0}$$

$$\ddot{\phi}_{1} = \frac{\phi_{1} - \phi_{0} - h\dot{\phi}_{0}}{\alpha h^{2}} - \frac{c_{1}\ddot{\phi}_{0}}{\alpha}$$

$$c_{1} = \frac{1}{2} - \alpha$$

$$c_{2} = 1 - \frac{\delta}{2\alpha}$$

#### **Two-Dimensional Wave Equations - Example 2**

**Newmark's Method** – Evaluating the differential equation at *t* = 0 yields

$$\begin{split} \mathbf{M} \ddot{\phi}_{0} &= \mathbf{F} - \mathbf{K} \phi_{0} \implies \ddot{\phi}_{0} = \mathbf{M}^{-1} \left( \mathbf{F} - \mathbf{K} \phi_{0}^{(0)} \right)^{(0)=0} \\ \ddot{\phi}_{0} &= \begin{bmatrix} 47.02041 & -11.75510 & 2.93878 & -11.75510 \\ -11.75510 & 23.51020 & -5.87755 & 2.93878 \\ 2.93878 & -5.87755 & 11.75510 & -5.87755 \\ -11.75510 & 2.93878 & -5.87755 & 23.51020 \end{bmatrix} \begin{bmatrix} 0.027778 \\ 0.013889 \\ 0.006944 \\ 0.013889 \end{bmatrix} \\ &= \begin{cases} 1.00000 \\ 0.00000 \\ 0.00000 \\ 0.00000 \\ 0.00000 \\ 0.00000 \end{bmatrix}$$

#### TIME-DEPENDENT PROBLEMS

#### **Two-Dimensional Wave Equations - Example 2**

**Newmark's Method** – Taking  $\alpha = 0.25$ ,  $\delta = 0.5$ , and h = 0.01 seconds yields:



#### **Two-Dimensional Wave Equations - Example 2**

**Newmark's Method –** The solution for  $\phi_1$  is

 $\phi_1 = \langle 4.996E-05 \ 1.284E-08 \ 6.595E-12 \ 1.284E-08 \rangle^T$ 

Now solve for the velocity at t = h

$$\dot{\phi}_{1} = \dot{\phi}_{0}^{\dot{\phi}_{0}(0)=0} + \frac{\delta(\phi_{1} - \phi_{0} - h\dot{\phi}_{0})}{\alpha h} + c_{2}h\ddot{\phi}_{0}^{\dot{\phi}_{0}(0)=0} + c_{2}h\ddot{\phi}_{0}^{\dot{\phi}_{0}(0)=0}$$

 $= \langle 9.9914E-03 \ 2.5670E-06 \ 1.3191E-09 \ 2.5670E-06 \rangle^{T}$ 

#### TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 2**

Newmark's Method – The equation for acceleration at t = 0 is  $\phi_0(0) = 0$  $\ddot{\phi}_1 = \frac{\phi_1 - h\dot{\phi}_0 - h\dot{\phi}_0}{\alpha h^2 h^2} - \frac{c_1\ddot{\phi}_0}{\alpha} = \frac{\phi_1}{\alpha h^2} - \frac{c_1\ddot{\phi}_0}{\alpha}$ 

Now solve for the acceleration at t = h

 $\ddot{\phi}_1 = \langle 0.99829 \quad 5.1341E-04 \quad 2.6381E-07 \quad 5.1341E-04 \rangle^T$ 

Two-Dimensional Wave Equations - Example 2 Newmark's Method – For n = 2:  $\phi_2 = \langle 1.9966E-04 \ 1.0259E-07 \ 7.9076E-11 \ 1.0259E-07 \rangle^T$   $\dot{\phi}_2 = \langle 0.01995 \ 1.5385E-05 \ 1.3177E-08 \ 1.5385E-05 \rangle^T$   $\ddot{\phi}_2 = \langle 0.99316 \ 2.0501E-03 \ 2.1078E-06 \ 2.0501E-03 \rangle^T$ For n = 3:  $\phi_3 = \langle 4.4859E-04 \ 4.2268E-07 \ 4.8051E-10 \ 4.2268E-07 \rangle^T$   $\dot{\phi}_3 = \langle 0.02984 \ 4.8633E-05 \ 6.7109E-08 \ 4.8633E-05 \rangle^T$  $\ddot{\phi}_3 = \langle 0.98464 \ 4.5996E-03 \ 8.6787E-06 \ 4.5996E-03 \rangle^T$ 

#### TIME-DEPENDENT PROBLEMS

Two-Dimensional Wave Equations - Example 2 Newmark's Method – For n = 11, t > 0.1 sec.



# **Two-Dimensional Wave Equations - Example 1**

The values of  $\phi_1$ ,  $\phi_2$ , and  $\phi_3$  for 0 < t < 10 sec. are show below:



## TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**

The velocity values of  $\phi_1$ ,  $\phi_2$ , and  $\phi_3$  for  $0 \le t \le 10$  sec. are:



# **Two-Dimensional Wave Equations - Example 1**

The acceleration values of  $\phi_1$ ,  $\phi_2$ , and  $\phi_3$  for  $0 \le t \le 10$  sec. are:



# TIME-DEPENDENT PROBLEMS

# **Two-Dimensional Wave Equations - Example 1**



## **Two-Dimensional Wave Equations - Example 1**



# TIME-DEPENDENT PROBLEMS

- **Closure** Time-dependent problems are inherently more difficult and expensive to solve than their corresponding steady-state counterparts.
- The expense of generating the global matrices is higher for the time-dependent problems because of the necessity of computing the mass matrices.
- The main extra expense, however, is in solving the resulting time-dependent global equations.

## **Two-Dimensional Wave Equations**

- **Closure -** For an analytical approach to the solution, additional expense is incurred in terms of having to determine eigenvalues and eigenvectors.
- The actual amount of expense depends on the specific form of the stiffness and mass matrices and the algorithm used, but in any case it is significantly in excess of the expense of solving the single set of linear algebraic equations associated with the steady-state problem.

# TIME-DEPENDENT PROBLEMS

- **Closure** For a time domain integration technique, the additional expense is clearly related to the number of time steps necessary to trace out the desired time history.
- In addition to several matrix multiplications and additions, *each step* can involve the solution of a set of linear algebraic equations.
- In some instances this expense can be minimized by using a decomposition that can be reused for the computation of the solution at each new time.

#### **Two-Dimensional Wave Equations**

- **Closure -** In this regard recall that the Euler and central difference algorithms require that the size of the time step not exceed a value proportional to the inverse of the largest eigenvalue.
- For large systems this critical step size can be very small resulting in many applications of the algorithm to trace out the time history.

# TIME-DEPENDENT PROBLEMS

- **Closure** The unconditionally stable Crank-Nicolson and Newmark algorithms, on the other hand, can be used with arbitrary step size that has been chosen so as to accurately integrate the lower modes, with significant improvement in the expense relative to the conditionally stable Euler and central difference algorithms.
- There are of course other algorithms available that are specifically tailored to address other numerical issues.

# End of Chapter 4d